

# **Quantum Time Series**

# Gábor Fáth

ELTE RiskLab & Department of Physics of Complex Systems

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## **Contributors**:

Zoltán Udvarnoki (ELTE)

Miklós Werner (BME, Wigner)

Örs Legeza (Wigner)

## Pricing and risk managing options



Figure 1.1: The S&P volatility surface as of June 20, 2013.

- Options modeled with stochastic volatility models
- Volatility is "rough" (Gatheral 2014)
- E.g.: Rough Fractional Stochastic Volatility (RFSV) model

$$dS_t = \mu_t S_t dt + \sigma_t S_t dZ_t$$
  

$$\sigma_t = \exp(X_t)$$
  

$$dX_t = \alpha (m - X_t) dt + \nu dB_t^H$$
 fOU process

- Fractional Brownian motion (fBm):  $B_t^H$ 
  - H = Hurst exponent
  - Self similarity (mono-fractal property):

 $\{B_H(at), t \in \mathbb{R}\} \stackrel{law}{=} \{a^H B_H(t), t \in \mathbb{R}\}$ 

Definition: **fBm** is a centered Gaussian process with auto-covariance

$$cov(B_H(s), B_H(t)) = \frac{var(B_H(1))}{2} \left( |t|^{2H} + |s|^{2H} - |t-s|^{2H} \right)$$

Integral representations:

• Mandelbrot-Van Ness

$$B_{H}(t) = B_{H}(0) + rac{1}{\Gamma(H+1/2)} \left\{ \int_{-\infty}^{0} \left[ (t-s)^{H-1/2} - (-s)^{H-1/2} 
ight] \, dB(s) + \int_{0}^{t} (t-s)^{H-1/2} \, dB(s) 
ight\}$$

• Molchan-Golosov

$$B_{H}(t) = \int_{0}^{t} K_{H}(t,s) \, dB(s) \qquad \qquad K_{H}(t,s) = rac{(t-s)^{H-rac{1}{2}}}{\Gamma(H+rac{1}{2})} \, _{2}F_{1}\left(H-rac{1}{2};\,rac{1}{2}-H;\,H+rac{1}{2};\,1-rac{t}{s}
ight).$$



## Cholesky:

- Using the known covariance structure of fBm
- Scaling  $O(n^3)$

## Circulant method (Davies-Harte)

- Simulate the process increments first (Fractional Gaussian noise, fGn)
- fBm sample follows from cumulative sum
- Covariance of fGn has "circulant" (Toeplitz) structure
- Can be diagonalized with Fourier transform
- Scaling  $O(n \ln(n))$

## Approximate methods

- Hybrid methods
- Kernel methods

$$\rho(k) = \mathbb{E}[\xi_1 \xi_{k+1}]$$
  
=  $\frac{1}{2n^{2H}} \left( |k+1|^{2H} - 2|k|^{2H} + |k-1|^{2H} \right)$ 

Covariance of the increments

## **Process of increments**

### Increments – Fractional Gaussian noise (fGn):

- Stationary
- Gaussian
- Autocorrelation (colored noise):

$$\rho(k) = \mathbb{E}[\xi_1 \xi_{k+1}]$$
  
=  $\frac{1}{2n^{2H}} \left( |k+1|^{2H} - 2|k|^{2H} + |k-1|^{2H} \right)$ 

• Asymptotically  $\rho(k) \sim 2H(2H-1) k^{2H-2} = \begin{cases} H < \frac{1}{2}: \text{ negative, fast decay, integrable} \\ H = \frac{1}{2}: \text{ zero, iid (White noise)} \\ H > \frac{1}{2}: \text{ positive, slow decay, non-integrable, long memory} \\ \text{Algebraic decay} \\ \text{exponent: } 2H-2 \end{cases}$ 

## Definition: Generalized Bernoulli Process

We will define stationary process,  $\{X_i, i \in \mathbb{N}\}$ , where each  $X_i$  takes one of two possible outcomes, 0 or 1, with  $P(X_i = 1) = p, P(X_i = 0) = 1 - p$ , and

 $cov(X_i, X_j) = c'|i-j|^{2H-2}, i \neq j,$ 

The integrated process is the Fractional Binomial Process

Define  $B_n = \sum_{i=1}^n X_i$ . It follows that  $E(B_n) = np$ , and as  $n \to \infty$ ,

$$Var(B_n) \sim \begin{cases} \left( p(1-p) + \frac{c'}{2H-1} \right) n & H \in (0, 1/2), \\ c'n \ln n & H = 1/2, \\ \frac{c'}{2H-1} |n|^{2H}, & H \in (1/2, 1). \end{cases}$$

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© Open Access Published by De Gruyter Open Access March 1, 2021 Generalized Bernoulli process with longrange dependence and fractional binomial distribution

#### Jeonghwa Lee 🖂

From the journal Dependence Modeling https://doi.org/10.1515/demo-2021-0100

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The **spin-1/2 XXZ spin chain** model is an example of a strongly correlated 1D quantum lattice system. The Hamiltonian is:

$$H = \sum_{j=1}^{N} \frac{1}{2} \left( S_{j}^{+} S_{j+1}^{-} + S_{j}^{-} S_{j+1}^{+} \right) + \Delta S_{j}^{z} S_{j+1}^{z}$$

Here  $S_i^+$ ,  $S_i^-$  and  $S_i^z$  are the spin operators at site *j*, and  $\Delta$  is the anisotropy parameter.

The phase diagram of the XXZ model exhibits different phases based on the value of  $\Delta$ :

- **Isotropic AF** ( $\Delta = 1$ ): Gapless. Critical behavior with algebraic decay of correlations
- Gapless Phases ( $-1 < \Delta < 1$ ): Gapless. Critical behavior with different correlation exponents (Luttinger liquid)
- Antiferromagnetic Phase ( $\Delta < 1$ ): Gapped. Long-range order with antiferromagnetic correlations (Néel order)
- **Ferromagnetic Phase (** $\Delta > 1$ **):** Long-range order with ferromagnetic correlations

## **Gapless phase**

Asymptotics of correlation functions ( $\Delta$ -dependent exponents):

$$C^{x}(n) = (-)^{n} \frac{A}{4n^{\eta}} \left( 1 - \frac{B}{n^{4/\eta - 4}} \right) - \frac{\tilde{A}}{4n^{\eta + 1/\eta}} \left( 1 + \frac{\tilde{B}}{n^{2/\eta - 2}} \right) + \dots$$
$$C^{z}(n) = -\frac{1}{4\pi^{2}\eta} \frac{1}{n^{2}} \left( 1 + \frac{\tilde{B}_{z}}{n^{4/\eta - 4}} \frac{4 - 3\eta}{2 - 2\eta} \right) + (-)^{n} \frac{A_{z}}{4n^{1/\eta}} \left( 1 - \frac{B_{z}}{n^{2/\eta - 2}} \right) + \dots$$

$$\Delta = -\cos(\pi\eta), \qquad 0 < \eta < 1$$



Long-distance asymptotics of spin-spin correlation functions for the XXZ spin chain

<u>Sergei Lukyanov</u> <sup>a b</sup>, <u>Véronique Terras</u> <sup>a c</sup> 🖂

	Classical time series	Quantum chain
Integrated process	Fractional Binomial Process	Spin domain magnetization
Increment process	Generalized Bernoulli Process	Spin-1/2 quantum chain
Correlations	Power law (in time)	Power law (in space)
Fractal characteristics	Hurst exponent	Correlation exponents
Probability measure	Bernoulli probs conditional on history	Quantum ground state implied
Sample	Process trajectory	Spin chain configuration
Sampling algorithm	Cholesky, Circulant, Kernel	Quantum sampling from MPS



A **Matrix product state** (**MPS**) is a quantum state of many particles (in N sites), written in the following form:

$$|\Psi
angle = \sum_{\{s\}} {
m Tr} \Big[ A_1^{(s_1)} A_2^{(s_2)} \cdots A_N^{(s_N)} \Big] |s_1 s_2 \dots s_N
angle,$$

where  $A_i^{(s_i)}$  are complex, square matrices of order  $\chi$  (this dimension is called local dimension). Indices  $s_i$  go over states in the computational basis. For qubits, it is  $s_i \in \{0, 1\}$ . For qudits (d-level systems), it is  $s_i \in \{0, 1, \dots, d-1\}$ .

## **MPS Sampling**

## Perfect sampling with unitary tensor networks

Andrew J. Ferris and Guifre Vidal Phys. Rev. B **85**, 165146 – Published 30 April 2012



1. Draw 1st spin:

$$\rho_{1} = \operatorname{Tr}_{2...N} |\Psi\rangle \langle \Psi|$$
$$P(s_{1}) = \langle s_{1} | \rho_{1} | s_{1} \rangle$$
$$\mathsf{Draw} \ s_{1}$$

2. Draw 2nd spin (conditioned on  $s_1$ ):

$$\rho_{2}(s_{1}) = \frac{1}{P(s_{1})} \operatorname{Tr}_{3...N} \langle s_{1} | \Psi \rangle \langle \Psi | s_{1} \rangle$$
$$P(s_{2} | s_{1}) = \langle s_{2} | \rho_{2}(s_{1}) | s_{2} \rangle$$
$$\operatorname{Draw} s_{2}$$

3. Draw 3rd spin (conditioned on  $s_1, s_2$ ):

$$\rho_3(s_1, s_2) = \frac{1}{P(s_1, s_2)} \operatorname{Tr}_{4\dots N} \langle s_1, s_2 | \Psi \rangle \langle \Psi | s_1, s_2 \rangle$$

etc...



# **MPS** sampling

This can be done effectively in the MPS formalism:

$$\circ = M = A^{\uparrow}|\uparrow\rangle + A^{\downarrow}|\downarrow\rangle$$

$$i = [] = A^{\uparrow} \circ A^{\uparrow}$$

$$i = [] = A^{\downarrow} \circ A^{\downarrow}$$

$$(A^{\uparrow} \circ A^{\uparrow} + A^{\downarrow} \circ A^{\downarrow}) [] = \lambda_{+}[]$$

$$\circ \quad 2 \cdot \chi \cdot \chi$$

$$j = \chi^{2} \cdot \chi^{2}$$

$$j = \chi^{2} \cdot \chi^{2}$$

$$P(\uparrow |history) = \frac{\rho_{\uparrow\uparrow}}{\rho_{\uparrow\uparrow} + \rho_{\downarrow\downarrow}}$$

$$P(\downarrow |history) = \frac{\rho_{\downarrow\downarrow}}{\rho_{\uparrow\uparrow} + \rho_{\downarrow\downarrow}}$$

MPS is a variational Ansatz, it minimizes the energy for a suitable *M*.



- H: Hamiltonian
- M: MPS tensor
- $O = M x M^*$  transfer matrix
- L: Left eigenvector of O
- R: Right eigenvector of O
- $\lambda_+$ : Leading eigenvalue

Density Matrix Renormalization Group (DMRG)

Alternatively:

- Convex optimization problem for the *M* tensor in a  $2 \cdot \chi \cdot \chi$  dimensional space
- Gradient descent can be applied on a platform where Automatic Differentiation is implemented (e.g. TensorFlow)





# Thank you